

Non-Hermitian exceptional boundary mode and its application to topological laser

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Collaborators: Yuto Ashida, Takahiro Sagawa

KS, Y. Ashida, and T. Sagawa, Nat. Commun. **11**, 5745 (2020).

KS, Y. Ashida, and T. Sagawa, Phys. Rev. B **105**, 235426 (2022).

Outline

Introduction

- Non-Hermitian topology

Exceptional edge mode

- Fundamental concept
- Numerical results on a toy model

Applications to topological lasers

- Toy models and numerical results
- Proposal of photonic realization

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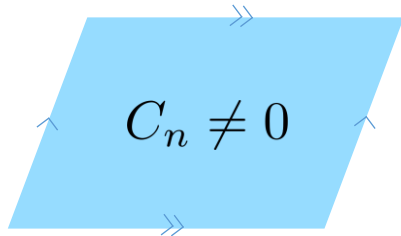
- Toy models and numerical results
- Proposal of photonic realization

Band Topology and Bulk-Edge Correspondence

Bulk (periodic boundary)

$$H(\mathbf{k})|\psi_n(\mathbf{k})\rangle = E_n(\mathbf{k})|\psi_n(\mathbf{k})\rangle$$

(Eigenequation of each wavenumber sector)

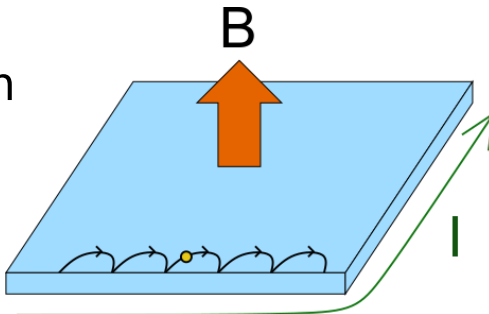


Chern number

$$C_n = \frac{1}{2\pi} \int_{\text{BZ}} \nabla_{\mathbf{k}} \times \mathbf{A}_n(\mathbf{k}) \cdot d\mathbf{S}$$

$$\mathbf{A}_n(\mathbf{k}) = i\langle\psi_n(\mathbf{k})|\nabla_{\mathbf{k}}|\psi_n(\mathbf{k})\rangle$$

e.g. Quantum Hall effect

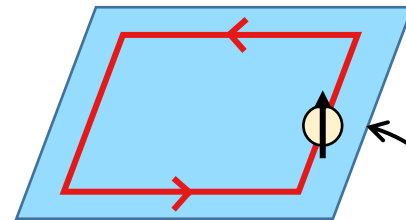


Bulk-edge correspondence



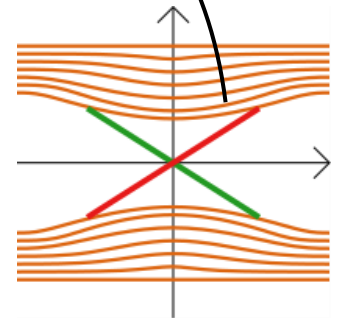
Edge (open boundary)

$$H|\Psi_n\rangle = E_n|\Psi_n\rangle$$



Edge modes

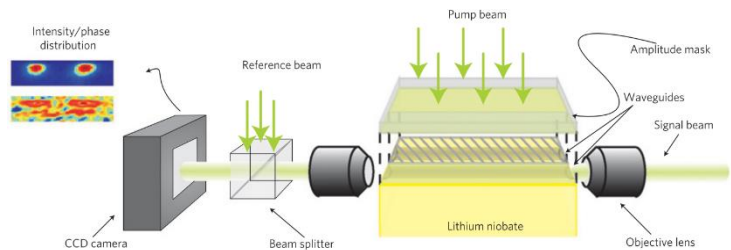
- **Localized** states
- Gapless dispersions
- **Backscattering-free** (unidirectional) current
- **Robustness** against disorders



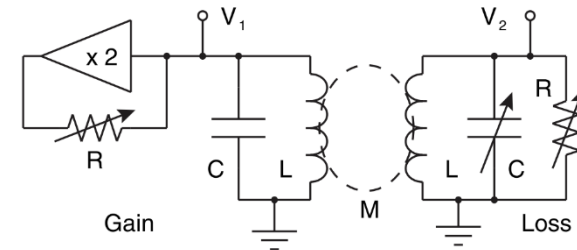
Topology of Non-Hermitian Hamiltonians

Non-Hermitian $H \neq H^\dagger$

- Description for systems **with gain-loss** and/or **nonreciprocity**
e.g. photonics, electrical circuits, cold atoms



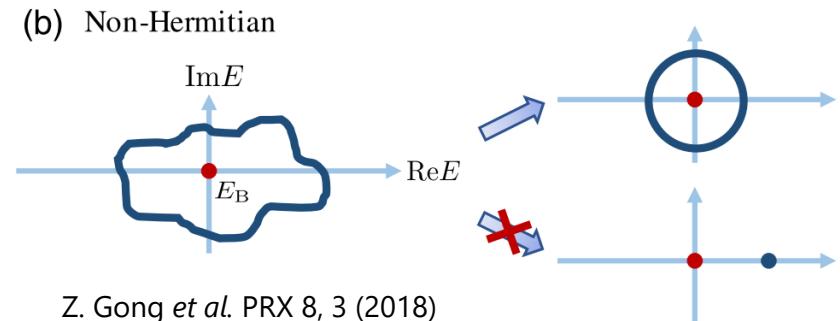
C. E. Rüter *et al.* Nat. Phys. 6, 192 (2010)



J. Schindler *et al.* Phys. Rev. A 84, 040101 (2011)

Unique topological features

- Novel classification utilizing complex eigenvalues
- Topological protection of **exceptional points**



Z. Gong *et al.* PRX 8, 3 (2018)

Exceptional Points: Non-Hermitian Topological Gapless Points

Exceptional points (EPs)

= **Nondiagonalizable gapless points**

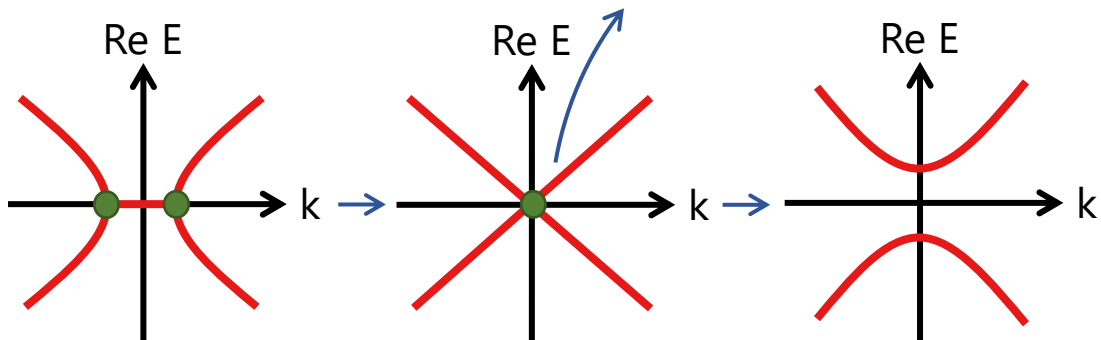
(cf. Jordan normal form)

cf) $H(k) = \begin{pmatrix} k & 1 \\ 0 & -k \end{pmatrix}$ exceptional point at $k = 0$

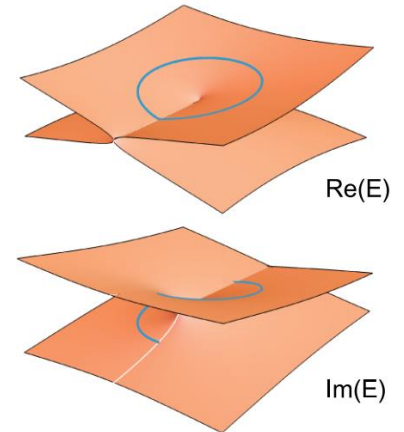


Topologically protected!!

→ Removed only via pair annihilations

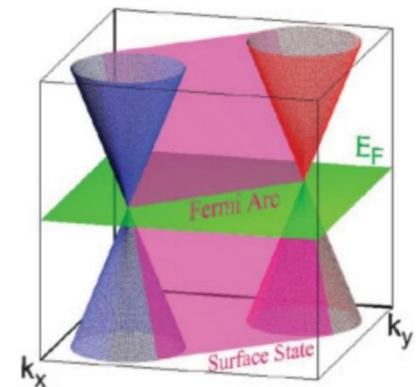


Green points: exceptional points



H. Shen *et al.* PRL 120, 146402 (2018)

Analogous to Weyl points



X. Wan *et al.*, PRB 83, 205101 (2011).

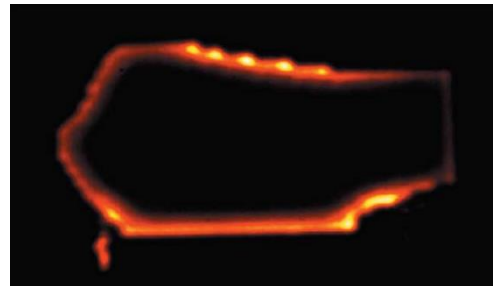
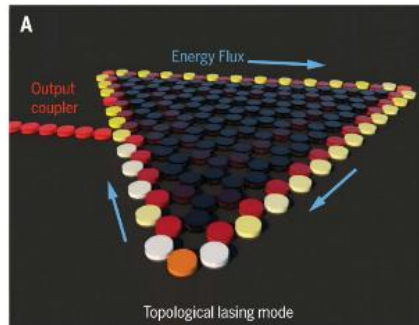
Topological Laser: Application of Non-Hermitian Topology

Topological laser

⇒ Amplifying boundary modes

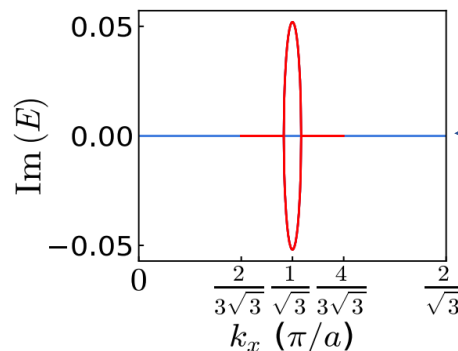
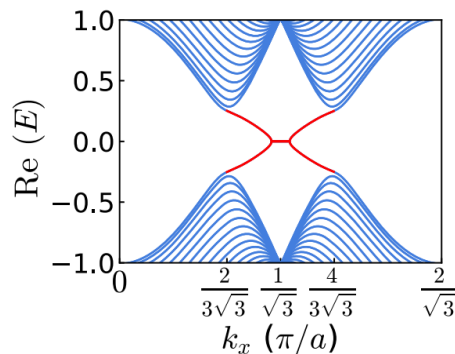
Advantage: robustness against disorders

- Topological insulators + **judicious gain** at the edge of the sample



- B. Bahari *et al.* Science 358, 636 (2017).
- G. Harari *et al.* Science 359, 1230 (2018).
- M. A. Bandres *et al.* Science 359, 1231 (2018).

- **Symmetry** protected lasing mode



PT symmetry

- K. Kawabata, K. Shiozaki, M. Ueda, M. Sato PRX 9, 041015 (2019).
- A. Y. Song *et al.* PRL 125, 033603 (2020).

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Applications to topological lasers

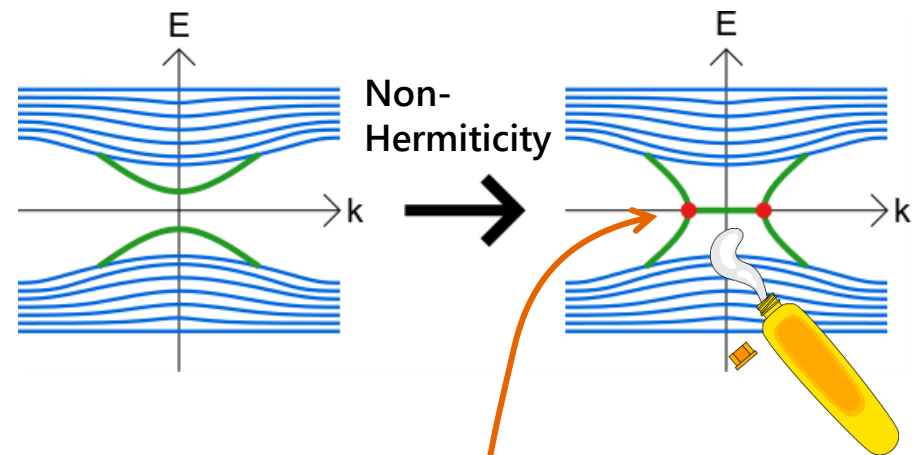
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Exceptional Edge Modes Independent of Bulk Topology

Exceptional edge modes (unique to non-Hermitian systems)

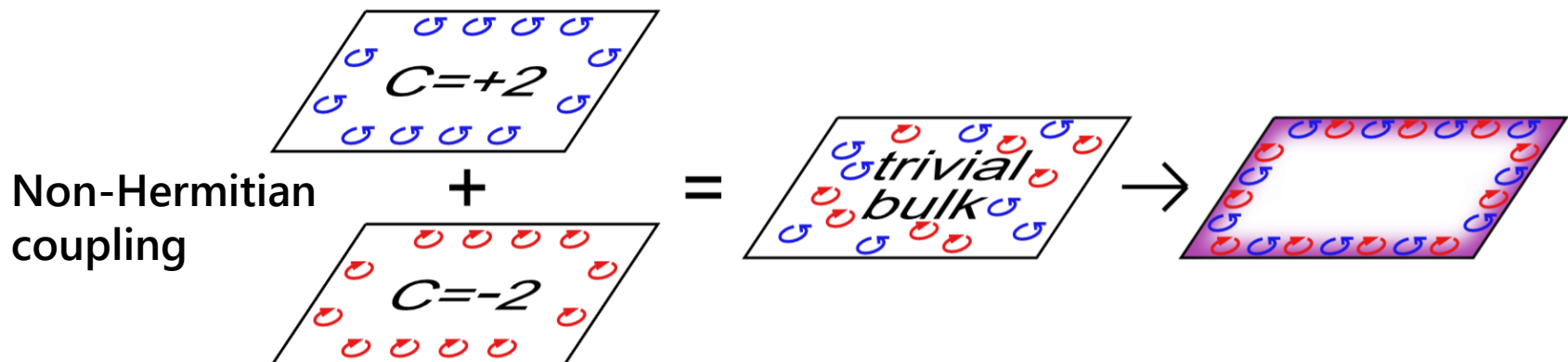
- Independent of the bulk topology
→ **Breakdown of the bulk-edge correspondence**

(different from skin effect)



Exceptional points protect edge modes like "glue"

Typical construction procedure



Toy Model for Exceptional Edge Modes

Toy model (tight-binding)

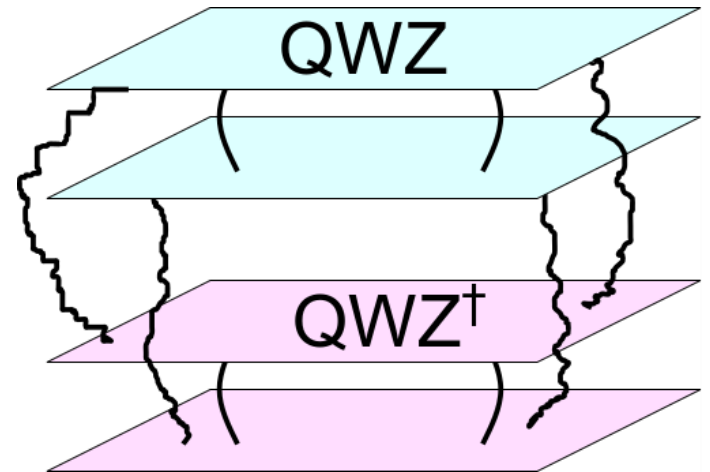
$$H = \begin{pmatrix} \boxed{H_{\text{QWZ}} & cI} & \boxed{ia\beta\sigma_x} & -i\alpha\sigma_x \\ cI & H_{\text{QWZ}} & i\alpha\sigma_x & \boxed{ia\beta'\sigma_x} \\ \boxed{ia\beta\sigma_x} & -i\alpha\sigma_x & \boxed{H_{\text{QWZ}}^*} & cI \\ i\alpha\sigma_x & \boxed{ia\beta'\sigma_x} & cI & H_{\text{QWZ}}^* \end{pmatrix}$$

Chern insulator

Non-Hermitian coupling

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

H_{QWZ} = 2×2 Hamiltonian of the QWZ model (Chern insulator)



Curve: Hermitian coupling

Wave curve: non-Hermitian coupling

Ref. QWZ model: X. L. Qi, Y. S. Wu, & S. C. Zhang PRB 74, 085308 (2006).

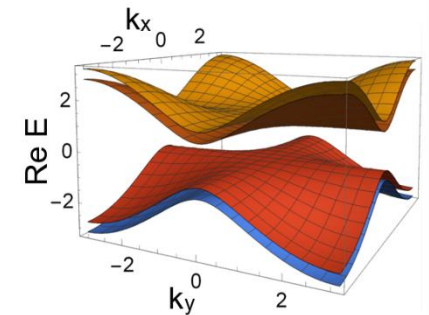
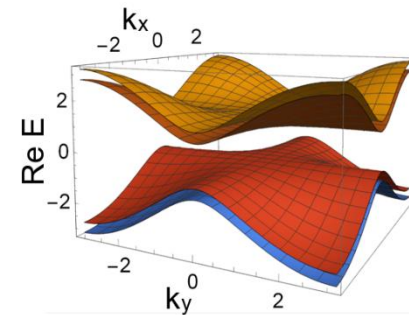
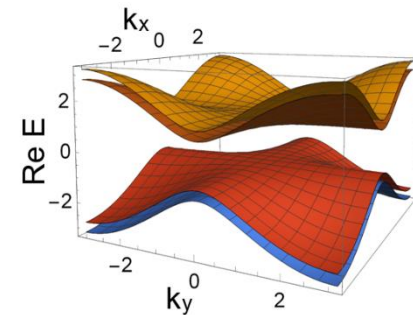
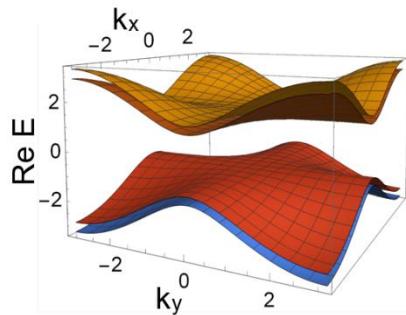
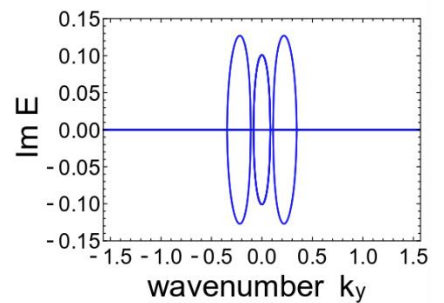
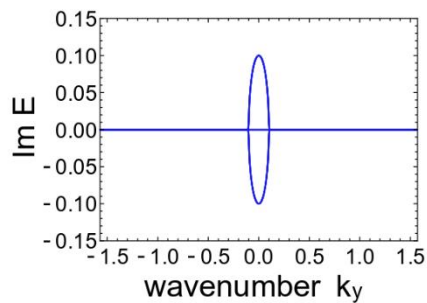
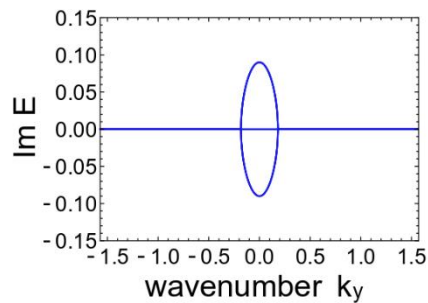
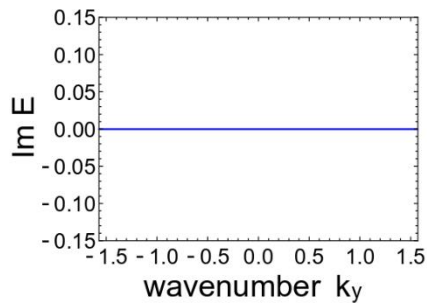
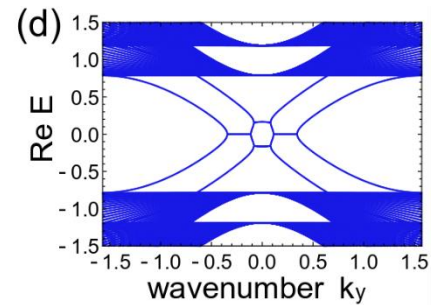
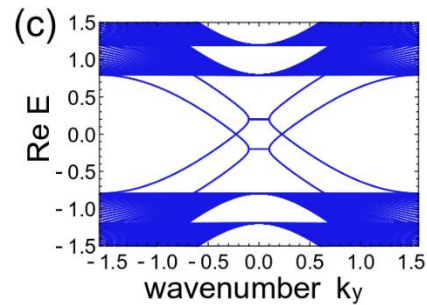
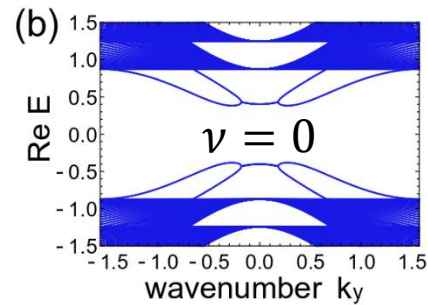
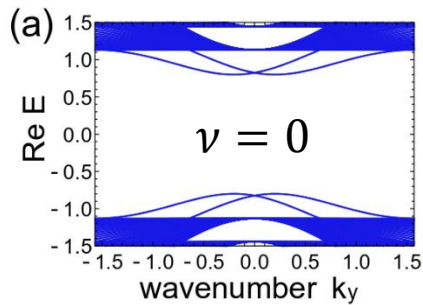
BHZ model: B. A. Bernevig, T. L. Hughes, & S. C. Zhang Science 314, 1757 (2006).

Independence of the Bulk Topology

0

Non-Hermiticity

large 



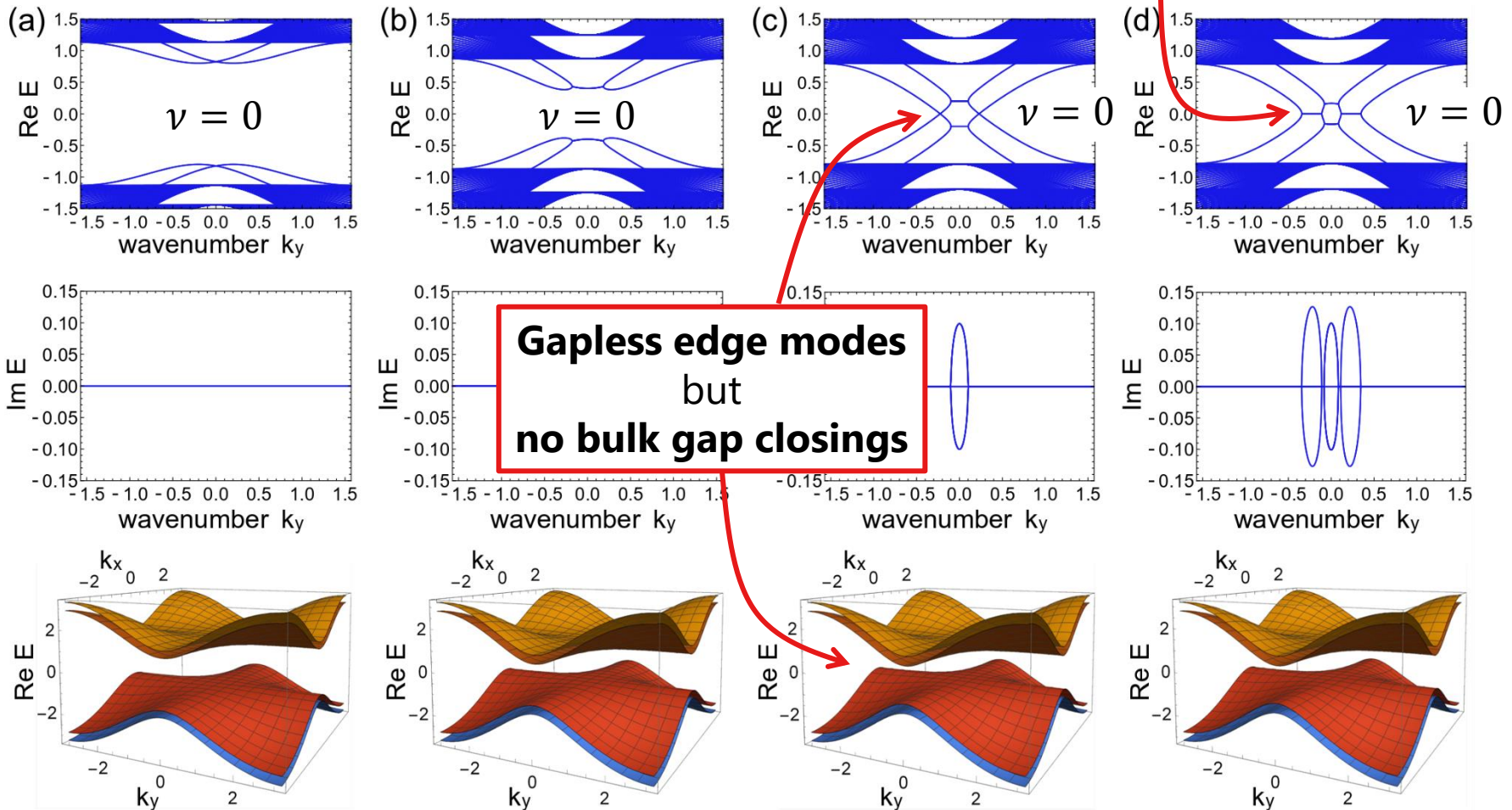
Independence of the Bulk Topology

0

Non-Hermiticity

large

Exceptional points



Topological Protection of Exceptional Points Under Symmetries

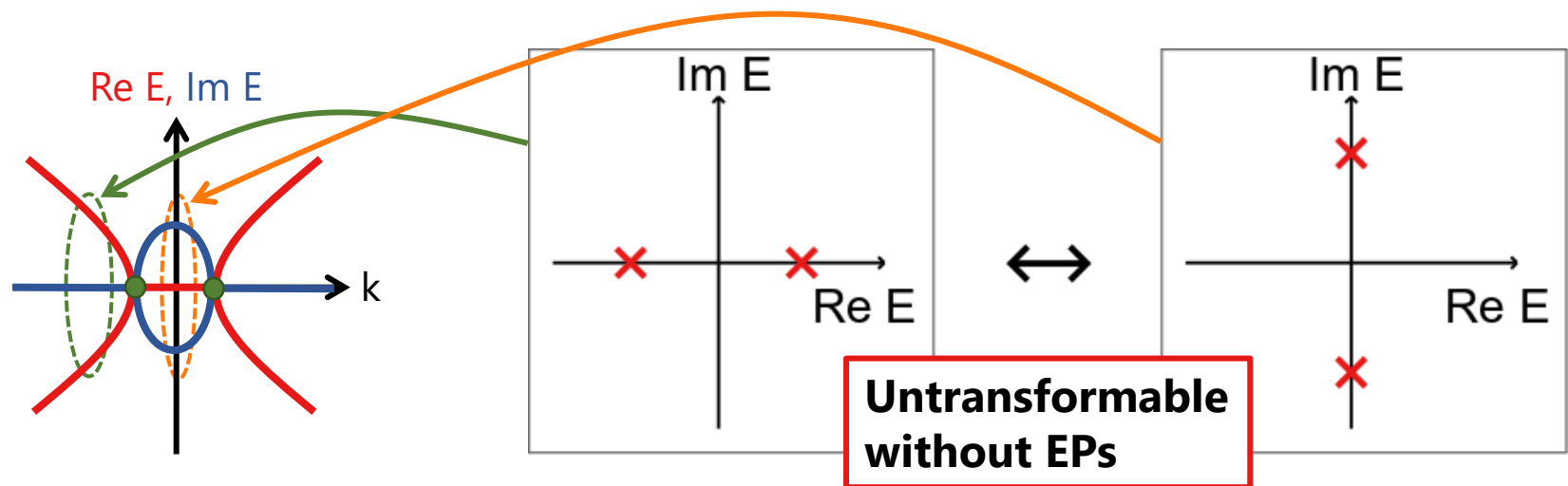
Exceptional points in 1D + symmetry

ex) PT (parity-time) symmetry

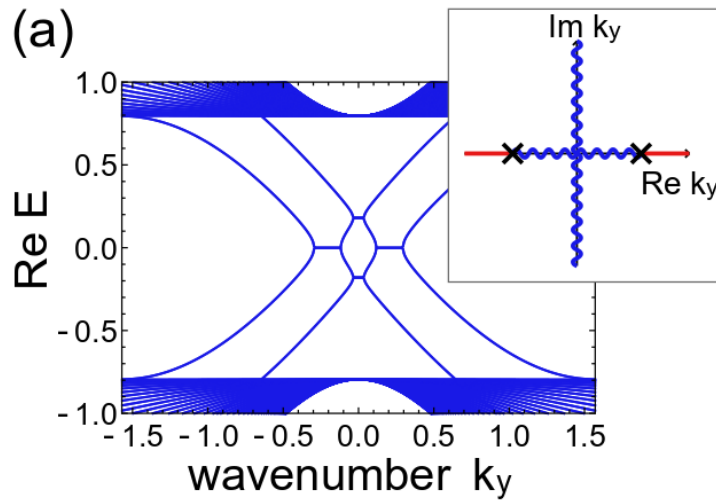
$$PTH(\mathbf{k})(PT)^{-1} = H^*(\mathbf{k})$$

→ Eigenvalues are real or pairs of complex conjugates.

→ EPs are protected!



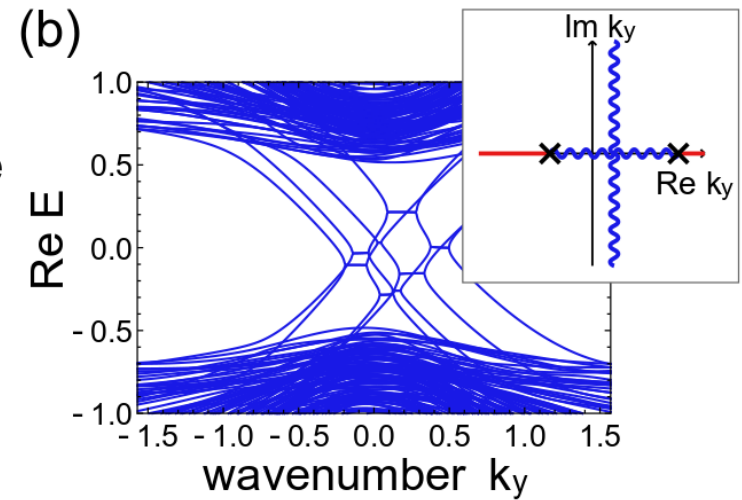
Robustness Against the Disorder



random
real on-site
potential

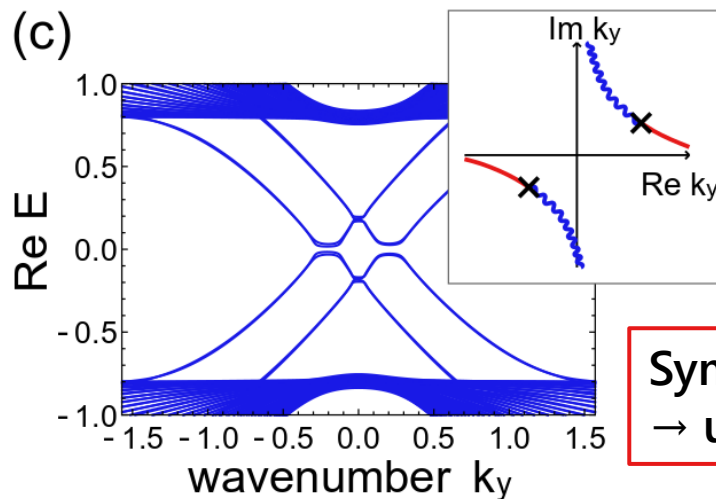


imaginary
noise in
coupling



PT-sym. preserved
→ robust

Consistent with the
symmetry protection
of EPs



Symmetries broken
→ unstable

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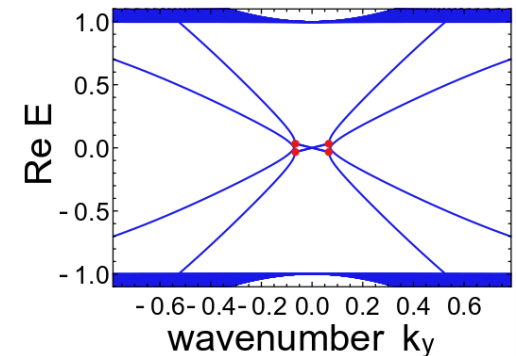
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Remained Problems on Topological Lasers

Q1. Propagation of lasing wave packets
without judicious gain

→ Without need of the knowledge
of boundary configurations

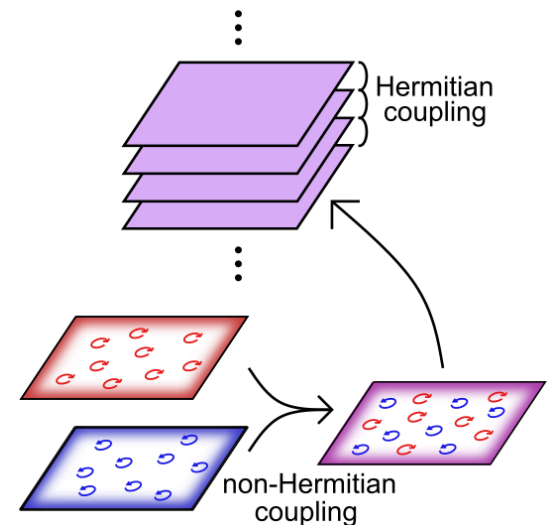
A1. Modified model of exceptional edge modes!



Q2. Topological lasers without
both judicious gain and symmetries

→ More robust topological lasers

**A2. Extension of exceptional edge modes
to 3D systems!**

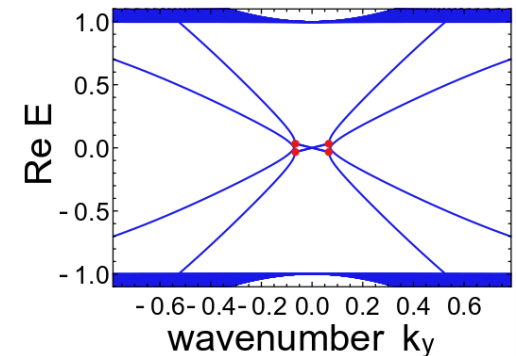


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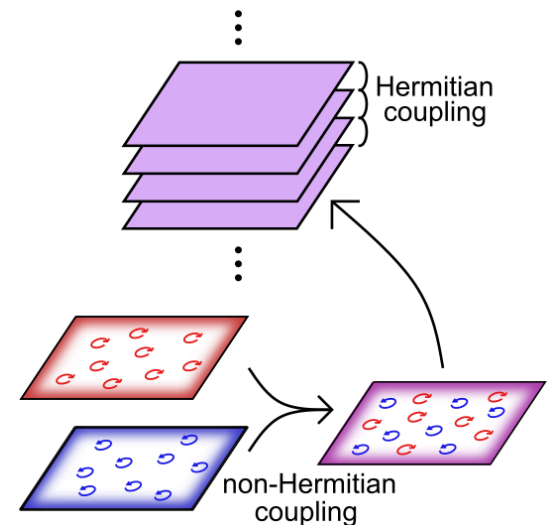
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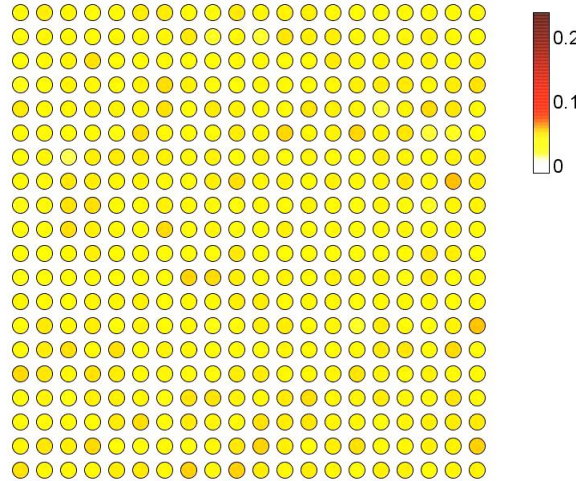


Application to 2D Topological Laser

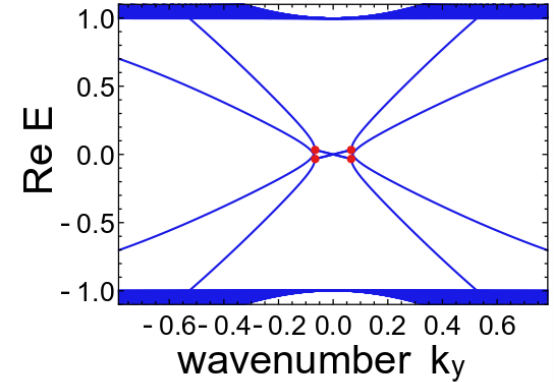
Hamiltonian

$$\begin{pmatrix} 2H_{\text{QWZ}} & i\beta\sigma_x \\ i\beta'\sigma_x & H_{\text{QWZ}}^* \end{pmatrix}$$

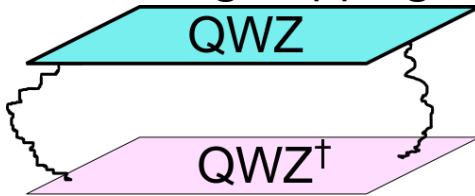
$$\beta, \beta' > 0$$



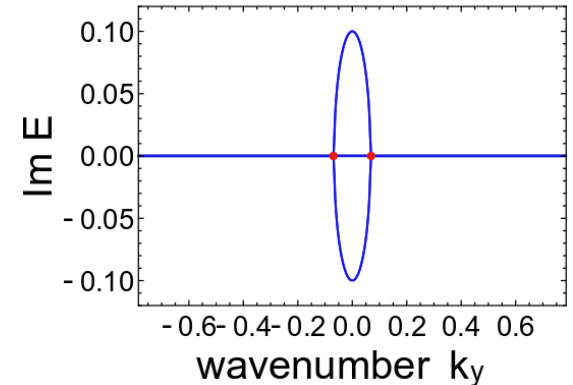
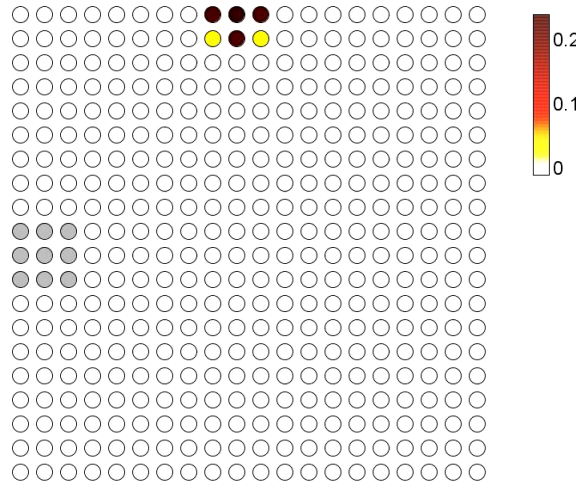
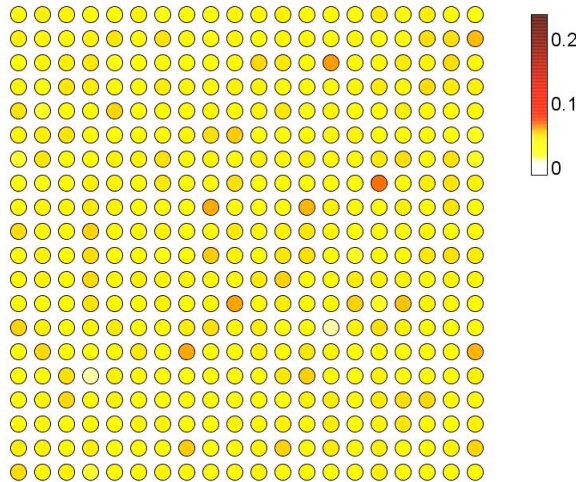
Edge band structure



Strong hopping



$$\beta = \beta' = 0$$



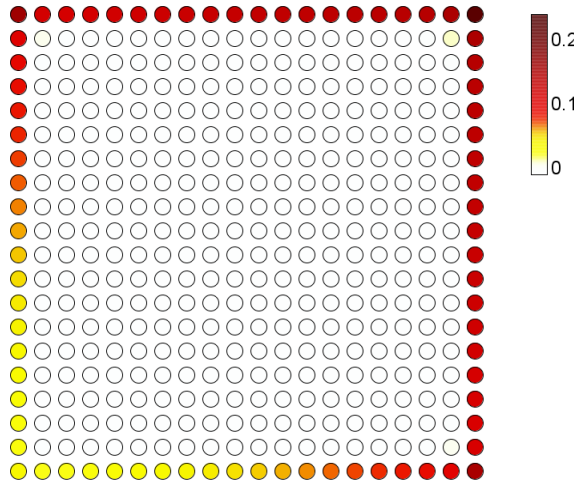
→ Satisfying prerequisites for a topological laser

Application to 2D Topological Laser

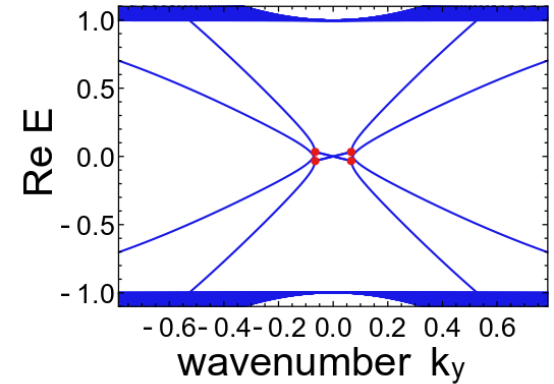
Hamiltonian

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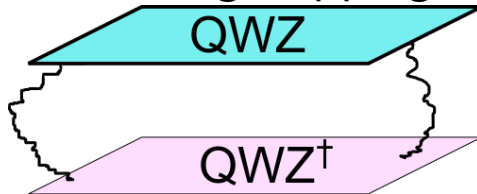
$$\beta, \beta' > 0$$



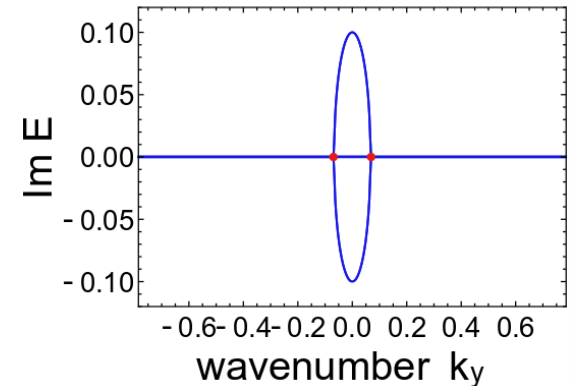
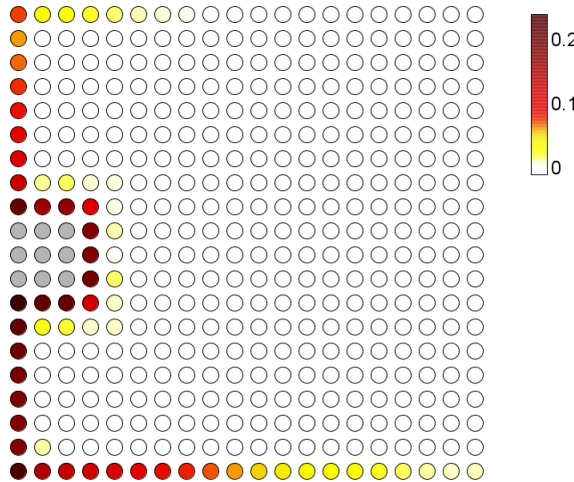
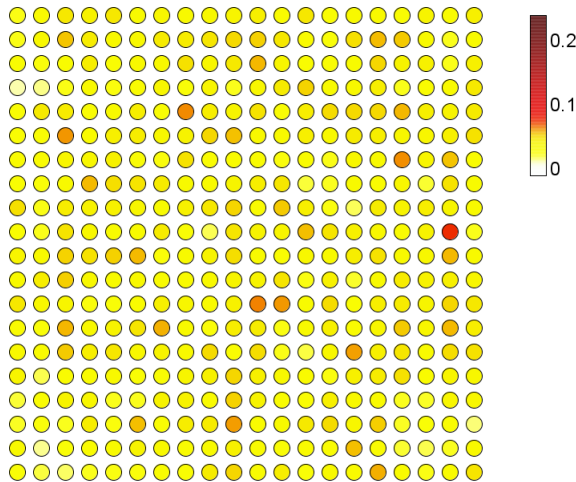
Edge band structure



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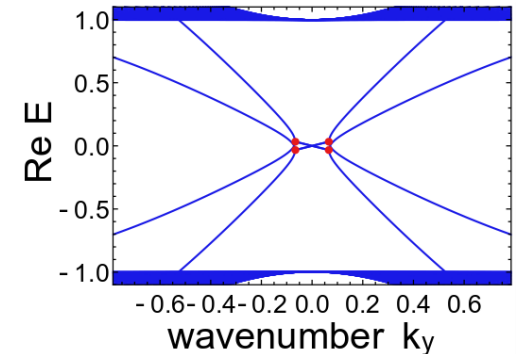
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→ Without need of the knowledge
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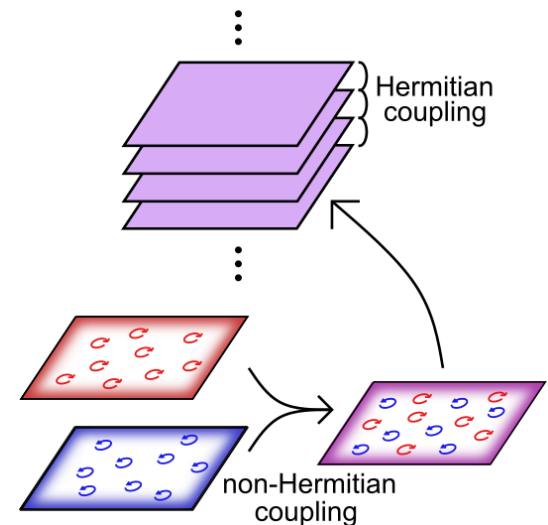
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Q2. Topological lasers without
both judicious gain and symmetries

→ More robust topological lasers

**A2. Extension of exceptional edge modes
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Construction Procedure of Three-Dimensional Topological Laser

No symmetry, 3D \rightarrow (Line-gap) topology is trivial.

(cf. K. Kawabata *et al.*, PRX 9, 041015 (2019).)



Exceptional points in boundary bands (2D) are stable.

\rightarrow extension of exceptional edge modes

Weak topological insulator

(accumulation of quantum spin Hall systems)

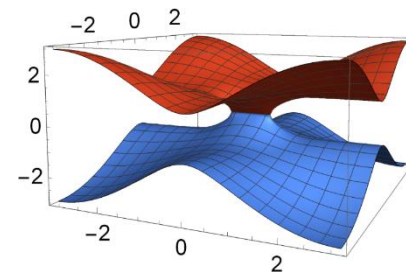
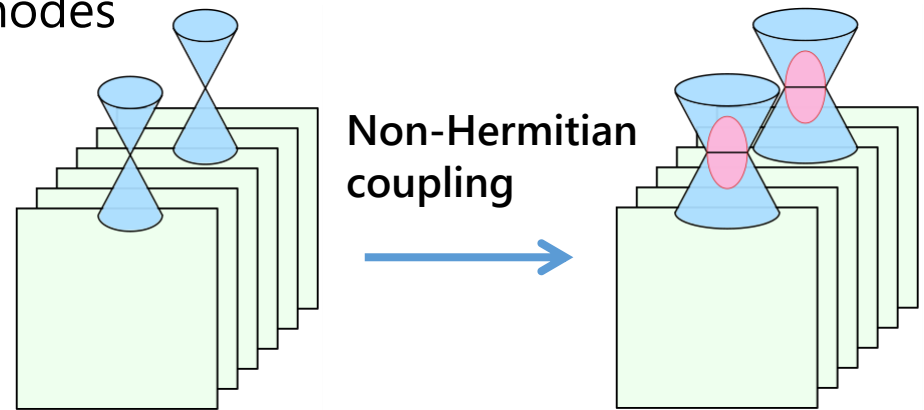
\rightarrow two Dirac cones

+

Non-Hermitian spin coupling

\rightarrow splitting into exceptional points

\rightarrow large imaginary parts of eigenvalues
(\Rightarrow amplification)



cf. One pair of EPs
 \rightarrow unstable under OBC

Toy Model: Weak Topological Insulator + Non-Hermitian Coupling

1. Chern insulator model

X. Qi, Y. S. Wu, S. C. Zhang PRB 74, 085308 (2006).

$$H(\mathbf{k}) = (u + \cos k_x + \cos k_y)\sigma_z + \sin k_x\sigma_x + \sin k_y\sigma_y$$

2. Coupling with its time-reversal counterpart (quantum spin Hall (QSH) system)

3. Non-Hermitian couplings

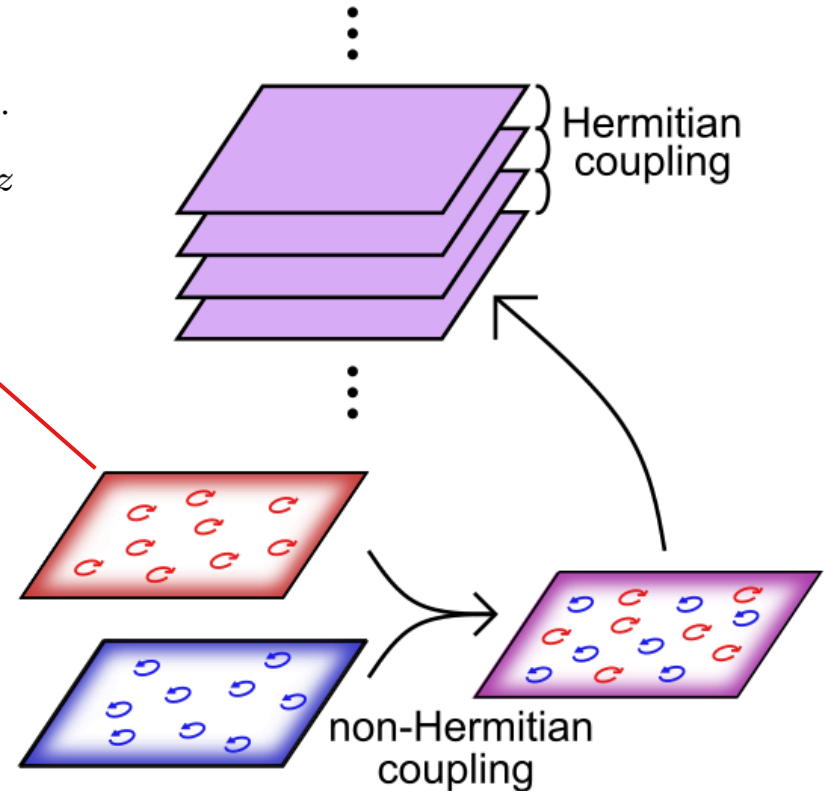
4. Accumulation

$$H(\mathbf{k}) = (u + \cos k_x + \cos k_y)I \otimes \sigma_z + \sin k_x\sigma_z \otimes \sigma_x + \sin k_y I \otimes \sigma_y + (c \sin k_z + i\gamma)\sigma_y \otimes \sigma_x$$

QSH system

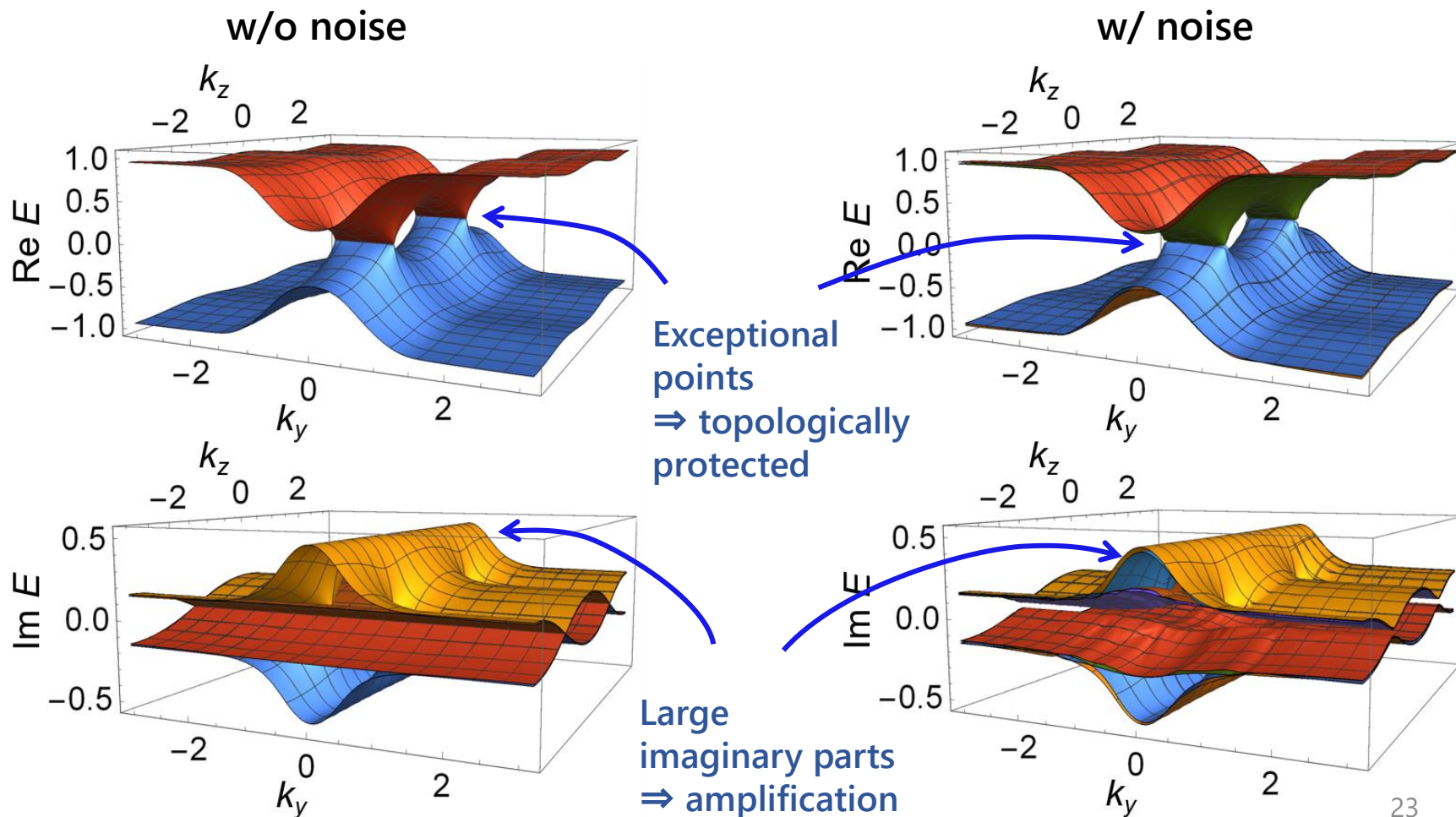
Interlayer coupling

Non-Hermitian coupling



Band Structure of the Surface Modes

$$H(\mathbf{k}) = (u + \cos k_x + \cos k_y)I \otimes \sigma_z + \sin k_x \sigma_z \otimes \sigma_x + \sin k_y I \otimes \sigma_y + (c \sin k_z + i\gamma)\sigma_y \otimes \sigma_x$$



Nonzero Group Velocity: Transfer of Lasing Wave Packets

Asymmetry between (effective) spins

→ **Nonzero group velocity** → **Transfer of lasing wave packets**

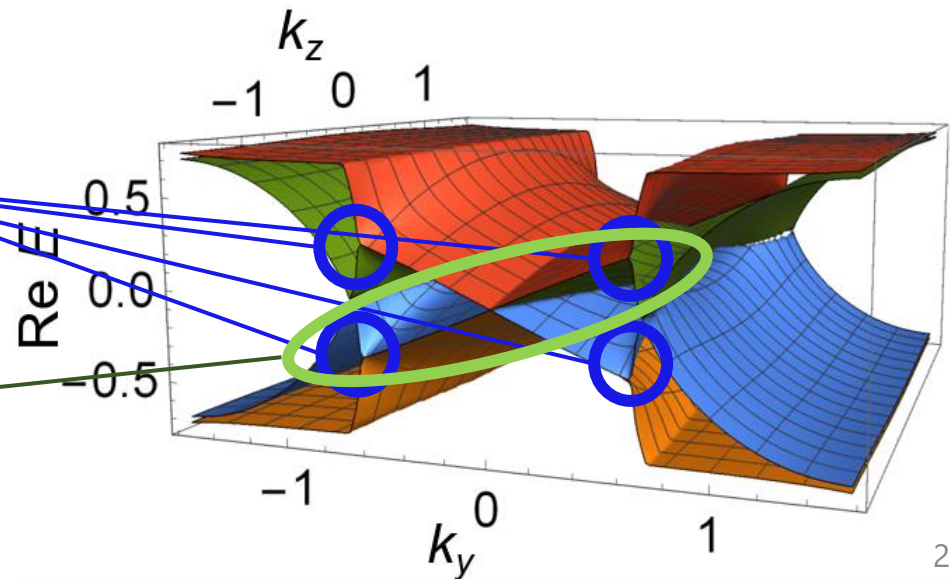
(cf. KS, Y. Ashida, and T. Sagawa, Nat. Commun. 11, 5745 (2020).)

$$\begin{aligned} H(\mathbf{k}) = & (u + \cos k_x + \cos k_y) (aI \otimes \sigma_z + \underline{b\sigma_z} \otimes \sigma_z) \\ & + \sin k_x (\underline{bI} \otimes \sigma_x + a\sigma_z \otimes \sigma_x) \\ & + \sin k_y (aI \otimes \sigma_y + \underline{b\sigma_z} \otimes \sigma_y) + (c \cos k_z + i\gamma)\sigma_y \otimes \sigma_x. \end{aligned}$$

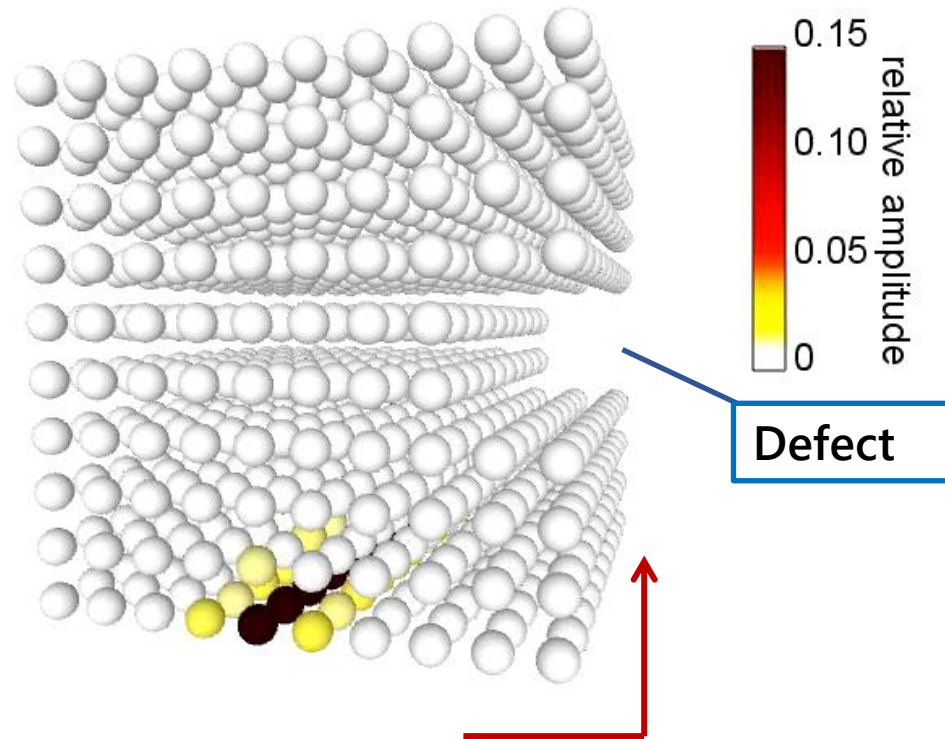
b : strength of asymmetry

Protection
by **exceptional points**

Boundary modes
with **nonzero group velocity**



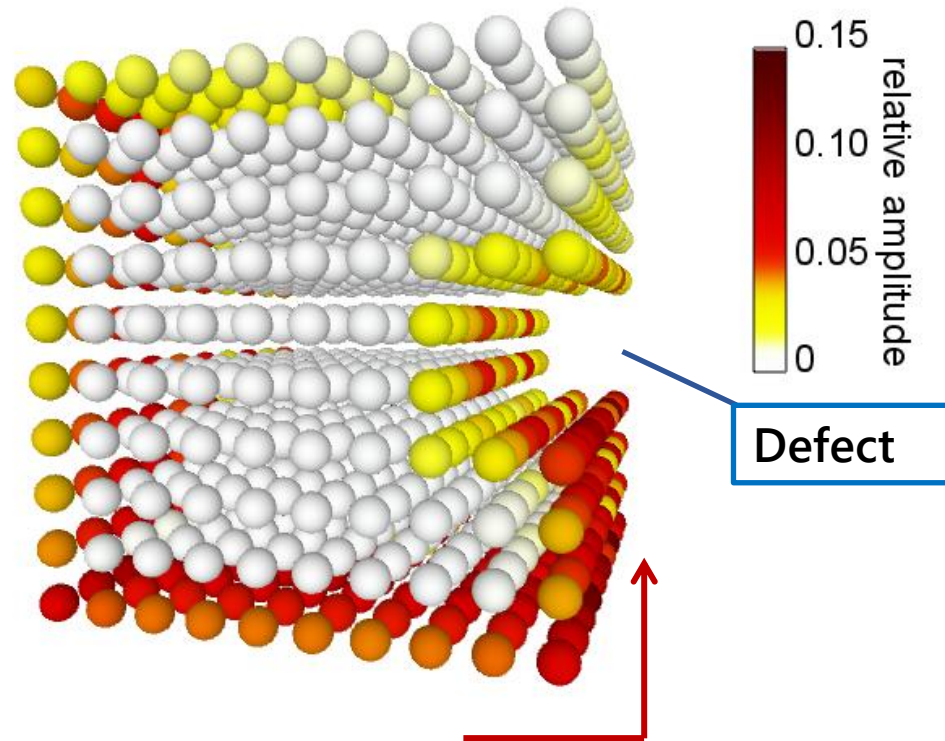
Real-Space Dynamics



Lasing wave packets along the edge

- Nonzero group velocity
 - Backscattering-free at the defect
- = **Robustness against disorders at the boundaries**

Real-Space Dynamics



Lasing wave packets along the edge

- Nonzero group velocity
- Scattering-free at the defect
= **Robustness against disorders at the boundaries**

Possible Optical Setup Using Ring-Resonator Arrays

Ring-resonator arrays

⇒ **quantum spin Hall systems**

- Effective spin: clockwise and counterclockwise modes
- Hopping: via evanescent light
- Artificial gauge field: phase-difference of propagating light

cf. M. Hafezi *et al.* Nat. Photonics 7, 1001 (2013).

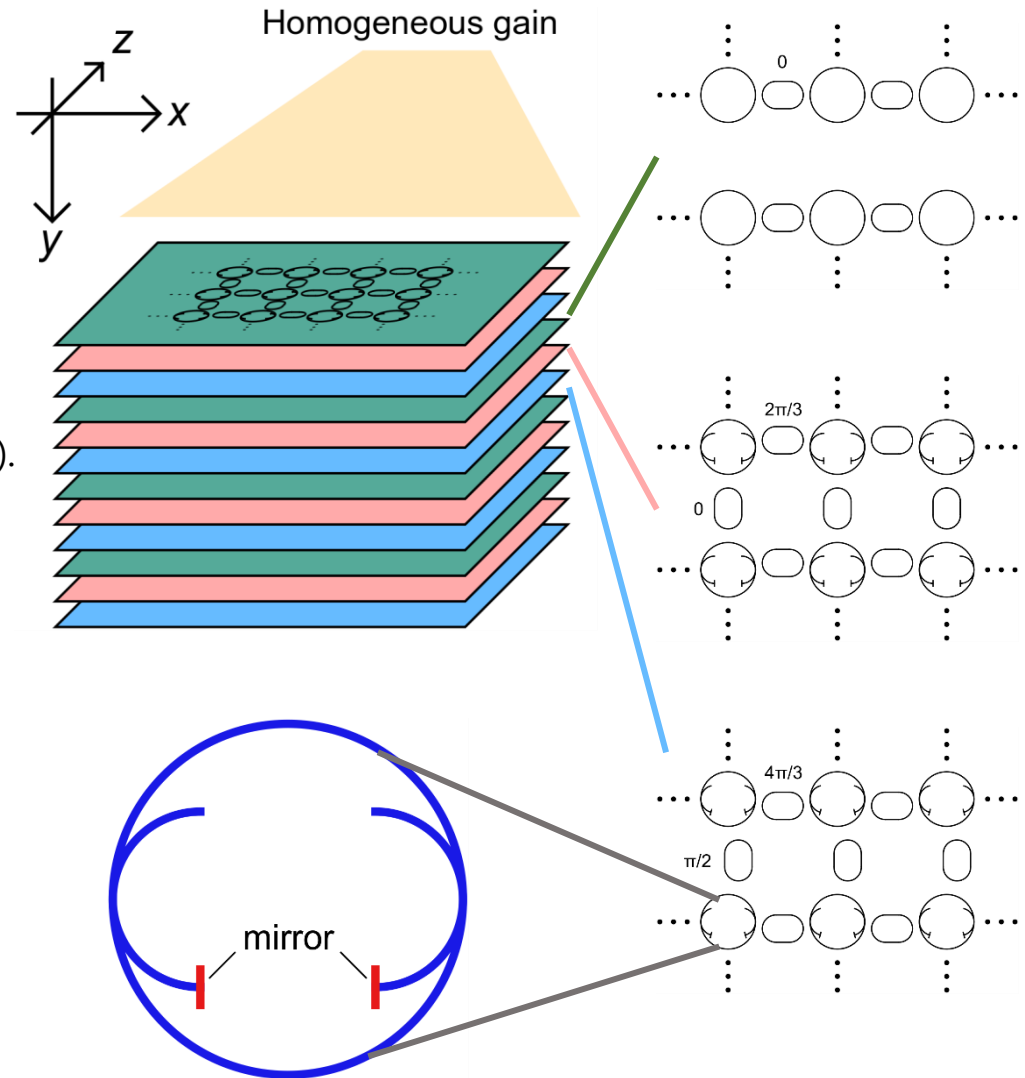
Accumulation

⇒ **Counterpart of weak topological insulator**

Waveguides with mirrors

⇒ **Non-Hermitian coupling**
(diminishing imbalance between CW and CCW modes)

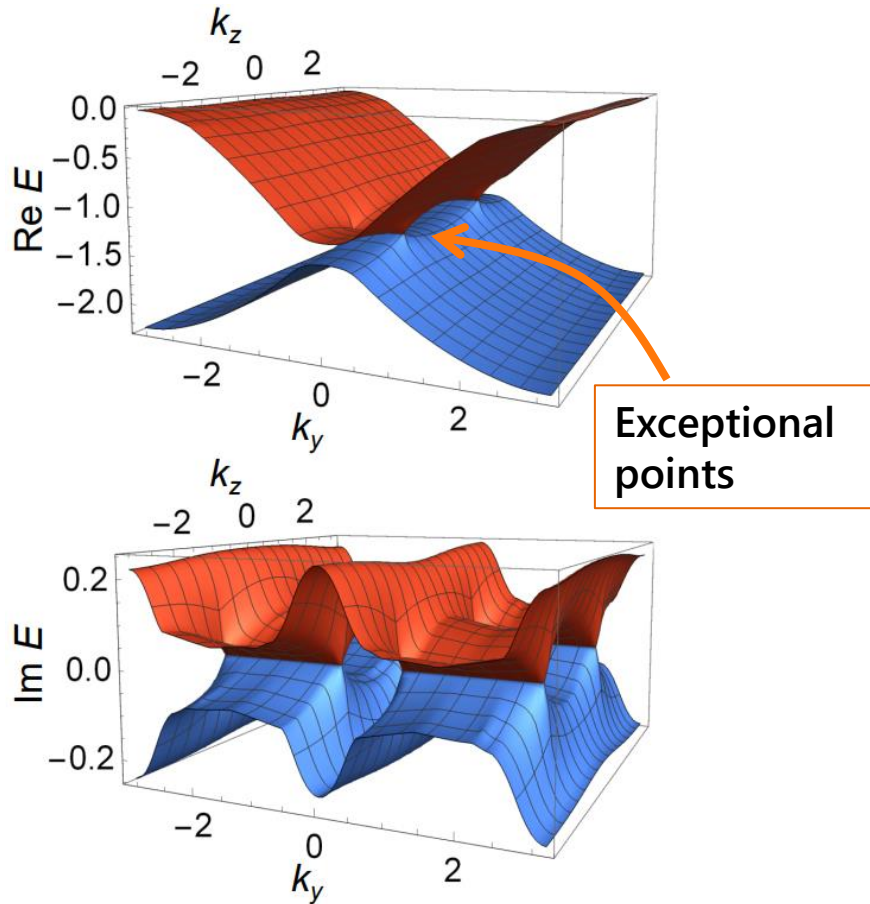
cf. A. Y. Song *et al.* PRL 125, 033603 (2020).



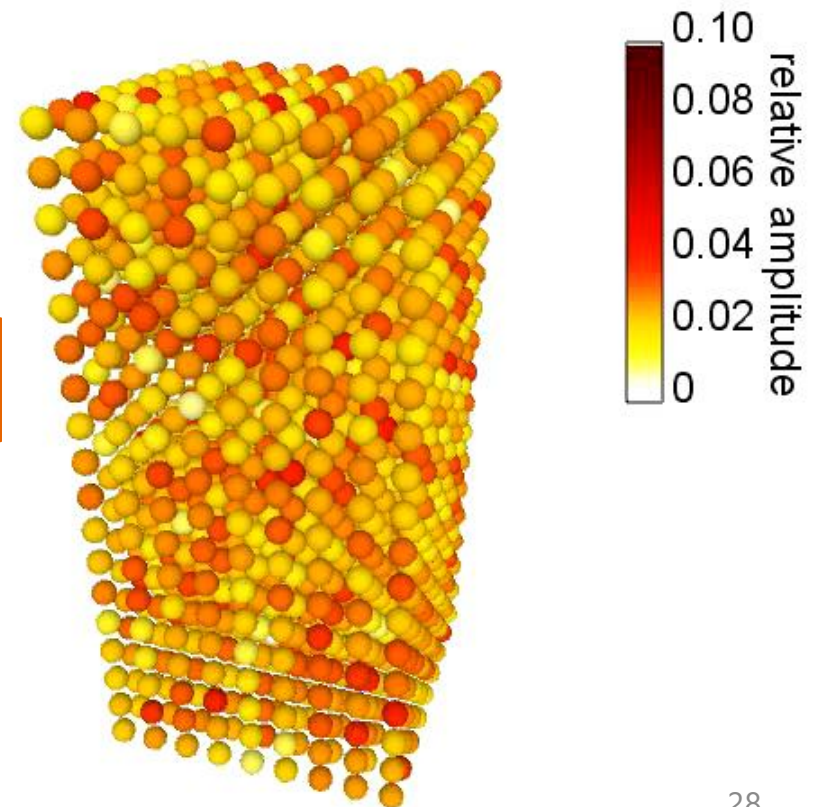
Numerical Results: Exceptional Surface Modes in Proposed Setup

$$H = \sum_{\mathbf{r}, s=\pm} (a_x e^{is\phi y} \hat{c}_{\mathbf{r},s}^\dagger \hat{c}_{\mathbf{r}+\mathbf{e}_x,s} + a_y \hat{c}_{\mathbf{r},s}^\dagger \hat{c}_{\mathbf{r}+\mathbf{e}_y,s}) + \text{H.c.} \\ + \sum_{\mathbf{r}} [a_z \sin(\phi y) (\hat{c}_{\mathbf{r},+}^\dagger \hat{c}_{\mathbf{r}+\mathbf{e}_z,-} + \hat{c}_{\mathbf{r},-}^\dagger \hat{c}_{\mathbf{r}+\mathbf{e}_z,+}) + \text{H.c.} + i\beta \sin(\phi y) (\hat{c}_{\mathbf{r},+}^\dagger \hat{c}_{\mathbf{r},-} + \hat{c}_{\mathbf{r},-}^\dagger \hat{c}_{\mathbf{r},+})]$$

Dispersion relation



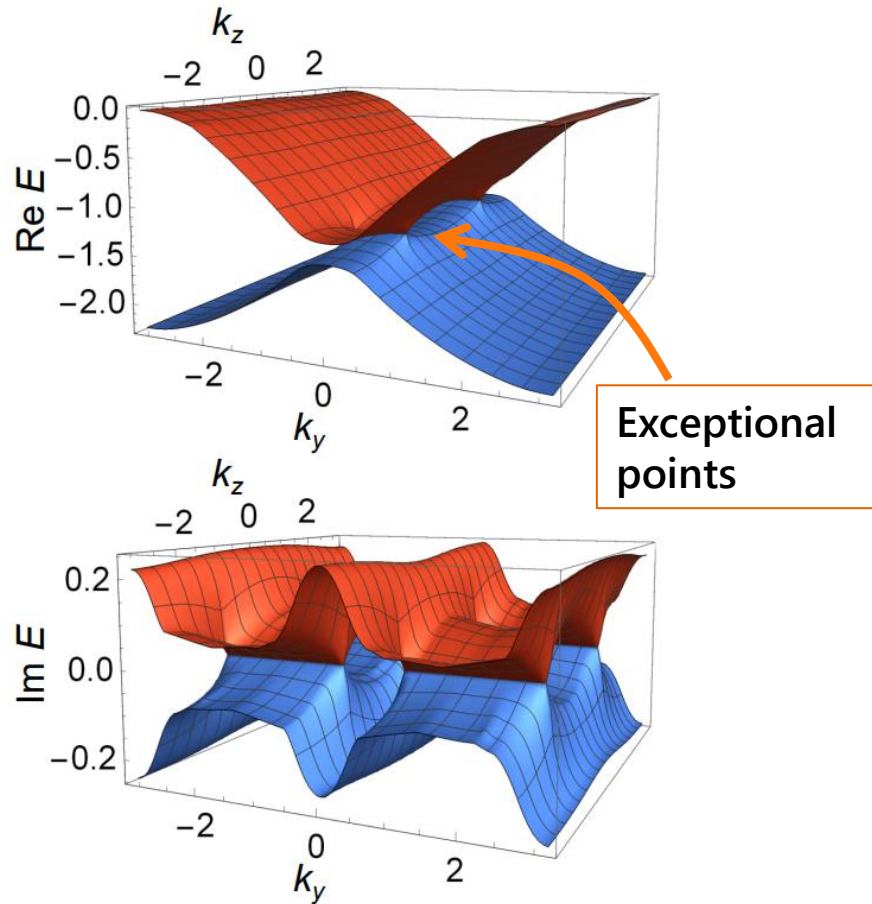
Real-space dynamics



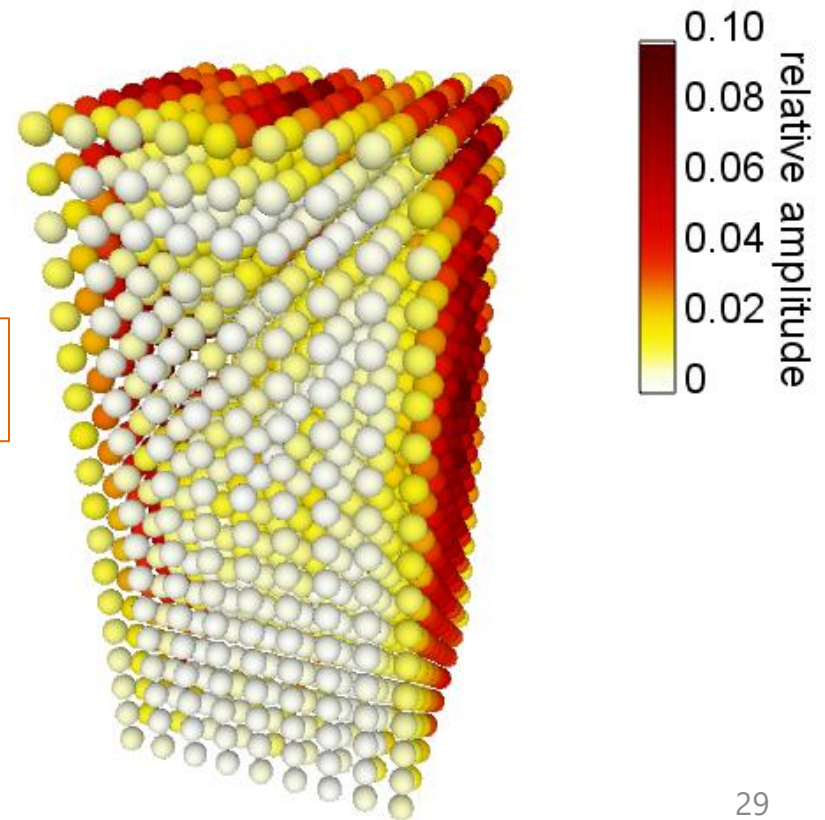
Numerical Results: Exceptional Surface Modes in Proposed Setup

$$\begin{aligned}
 H = & \sum_{\mathbf{r}, s=\pm} (a_x e^{is\phi y} \hat{c}_{\mathbf{r},s}^\dagger \hat{c}_{\mathbf{r}+\mathbf{e}_x,s} + a_y \hat{c}_{\mathbf{r},s}^\dagger \hat{c}_{\mathbf{r}+\mathbf{e}_y,s}) + \text{H.c.} \\
 & + \sum_{\mathbf{r}} [a_z \sin(\phi y) (\hat{c}_{\mathbf{r},+}^\dagger \hat{c}_{\mathbf{r}+\mathbf{e}_z,-} + \hat{c}_{\mathbf{r},-}^\dagger \hat{c}_{\mathbf{r}+\mathbf{e}_z,+}) + \text{H.c.} + i\beta \sin(\phi y) (\hat{c}_{\mathbf{r},+}^\dagger \hat{c}_{\mathbf{r},-} + \hat{c}_{\mathbf{r},-}^\dagger \hat{c}_{\mathbf{r},+})]
 \end{aligned}$$

Dispersion relation



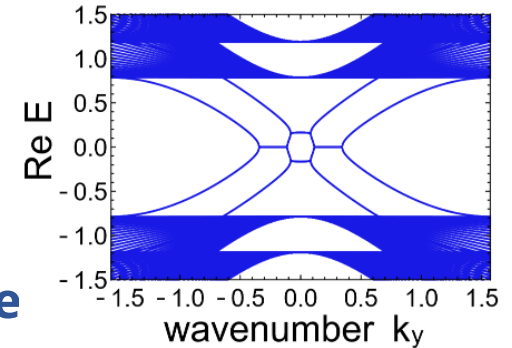
Real-space dynamics



Summary

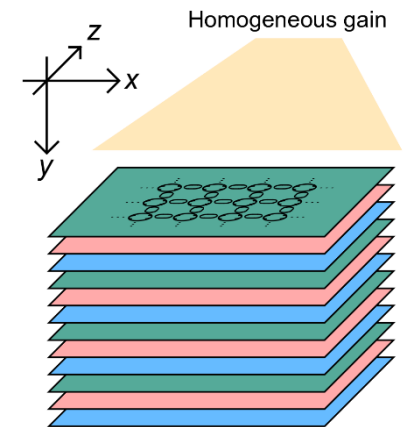
Exceptional edge modes

- Protected by **exceptional points**
- Robust even with a topologically trivial bulk
 - **Breakdown of the bulk-edge correspondence**



Applications to topological laser

- **Without judicious gains**
- In 3D, **no needs of symmetry** protections
- Proposal of photonic systems



Refs: KS, Y. Ashida, and T. Sagawa, Nat. Commun. **11**, 5745 (2020).
KS, Y. Ashida, and T. Sagawa, Phys. Rev. B 105, 235426 (2022).